

# CORRELATING MOLECULAR MODELLING AND EXPERIMENTAL DIFFUSIVITIES IN ZEOLITES

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# OUTLINE

1. Introduction.
2. Models and Methods.
3. Self-diffusion of methane – *n*-butane mixtures in silicalite-1.
4. Transport diffusivity of N<sub>2</sub> and CO<sub>2</sub> in silicalite-1.
5. *n*-alkanes in silicalite-1: Chain length dependence of self-diffusivity and activation energy.
6. Computer reconstruction of beds of zeolite crystals and simulation of long-range diffusion therein.
6. Summary

# ZEOLITES

## STRUCTURE

- framework geometry
- framework charge distribution
- size, shape, flexibility, and charge distribution of sorbed molecules

## MATERIAL PROPERTIES

### Sorption thermodynamics

- Henry's law constants
- Sorption isotherms

### Intracrystalline transport

- Self-diffusivities
- Transport diffusivities

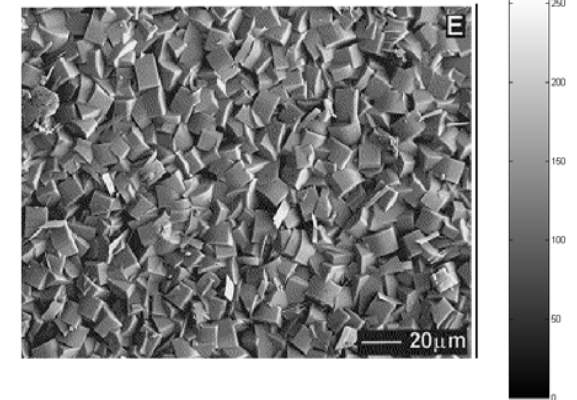
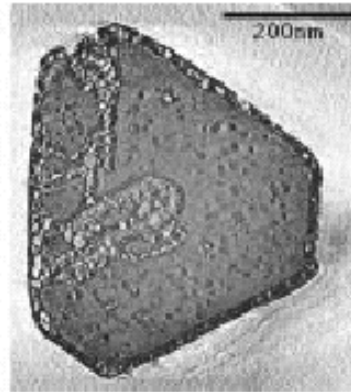
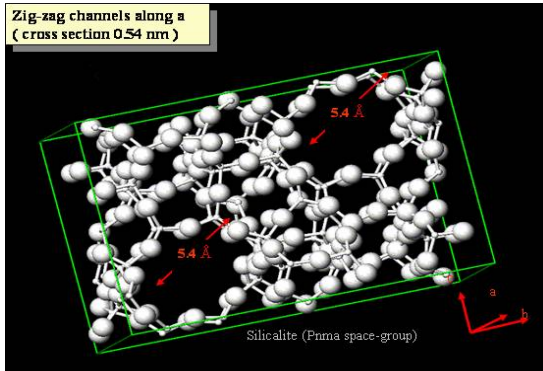
## MICROSCOPIC ASPECTS

- Siting of sorbate molecules
- Conformational changes upon sorption
- Modes and characteristic times of sorbate motion

## PERFORMANCE

Selectivity and rates in sorption separation and catalytic applications

# STRUCTURE AND TRANSPORT IN ZEOLITE MATERIALS MAY BE GOVERNED BY WIDE RANGE OF LENGTH SCALES



0.1 - 10 nm:  
Configurational diffusion  
in zeolite nanopores

10 nm – 1 μm:  
Defects, pore blocking in  
zeolite crystals

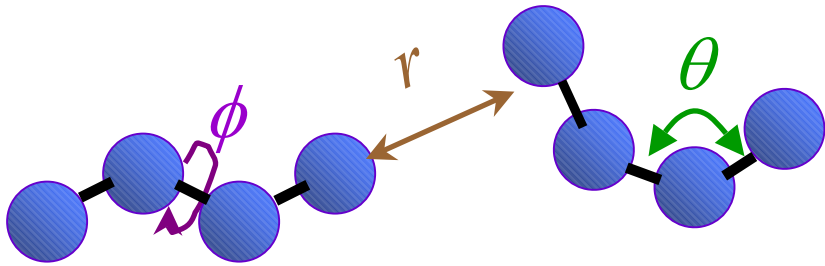
1 μm – 1 m: Molecular,  
Knudsen, surface diffusion  
in particles and beds

Prediction of performance in applications requires **hierarchical modelling approaches** that can address a wide range of length and time scales.

# MOLECULAR MODELLING

**Sorbates:** Use “molecular mechanics” force fields that give good predictions for fluid-phase properties

“united atom” representation for alkanes:



Lennard-Jones (dispersion attraction and excluded volume interactions):

$$V_{\text{LJ}}(r) = 4\epsilon \left[ \left( \frac{\sigma}{r} \right)^{12} - \left( \frac{\sigma}{r} \right)^6 \right]$$

Polar molecules (e.g., CO<sub>2</sub>):

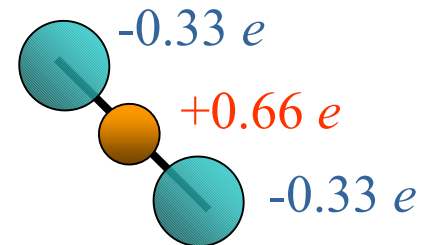
In addition, Coulomb interactions between partial charges:  $V_{\text{C}}(r) = \frac{z_i z_j e^2}{4\pi\epsilon_0 r}$

Bond angle bending:

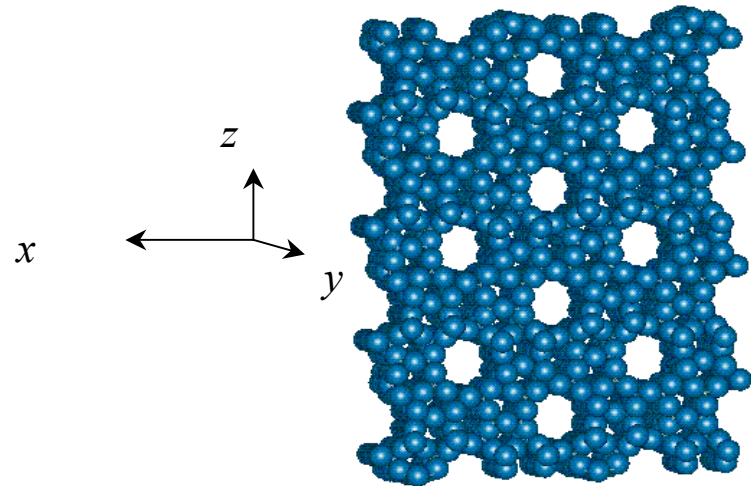
$$V_{\theta}(\theta) = \frac{1}{2} k_{\theta} (\theta - \theta_0)^2$$

Dihedral angle torsion:

$$V_{\phi}(\phi) = \sum_{k=0}^5 c_k \cos^k(\phi)$$



# MOLECULAR MODELS



## Silicalite-1



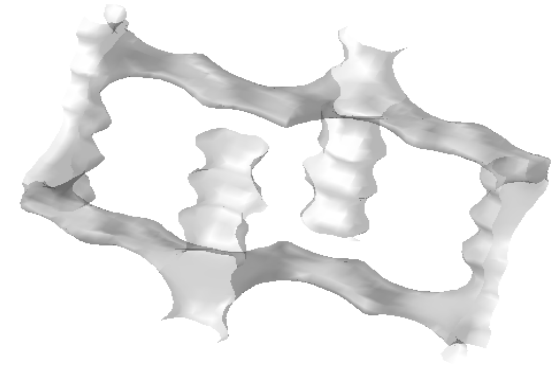
Pnma

$$a = 20.07 \text{ \AA}$$

$$b = 19.91 \text{ \AA}$$

$$c = 13.42 \text{ \AA}$$

(2 × 2 × 4  
unit cells shown)



**Zeolites:** Represented as sets of framework atoms and counterions at their crystallographically known positions.

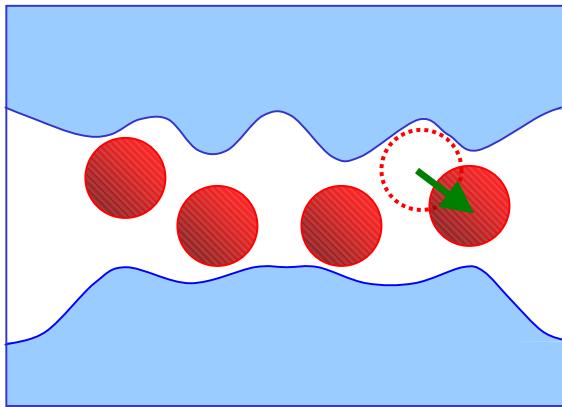
**Lennard-Jones** and **Coulomb** interactions with sites on sorbate molecules. Partial charges on framework from DFT calculations.

Rigid framework model allows **pretabulation** of the zeolite field experienced by sorbate sites: 100 fold savings in CPU time.

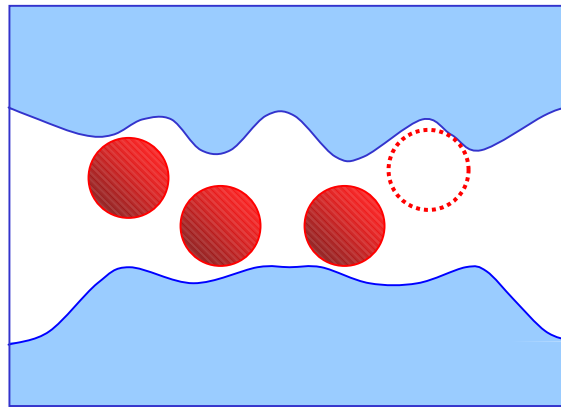
# PREDICTION OF SORPTION ISOTHERMS

Grand Canonical Monte Carlo (GCMC) Simulation: Sorbate-zeolite system simulated under constant  $V_s, T, f$  ← Sorbate fugacity

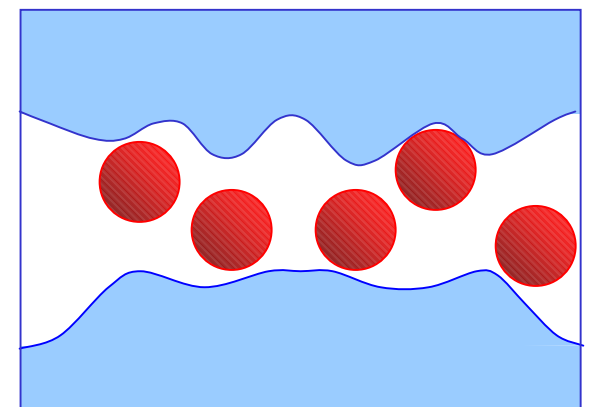
Elementary moves attempted:



Sorbate molecule  
displacement



Sorbate molecule  
removal



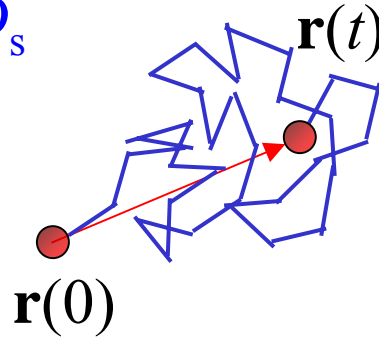
Sorbate molecule  
addition

Move acceptance criteria designed so that the sequence of configurations generated samples the equilibrium probability distribution of the grand canonical ensemble at prescribed  $V_s, T, f$ .

$$\langle N \rangle = f(V_s, T, f) \rightarrow \text{Equilibrium adsorption isotherm at } T$$

# PREDICTION OF DIFFUSIVITY: Molecular Dynamics (MD) Simulations

Self-Diffusivity  $D_s$



$$D_s = \lim_{t \rightarrow \infty} \frac{\langle [\mathbf{r}(t) - \mathbf{r}(0)]^2 \rangle}{6t}$$

(Einstein, 1905)

Transport Diffusivity  $D_t$

$$\mathbf{J} = -D_t \nabla c$$

flux,  $\text{mol m}^{-2} \text{s}^{-1}$

concentration,  
 $\text{mol m}^{-3}$  crystal

Corrected Diffusivity  $D_o$

L.S.Darken, *Trans. AIME* **1948**, 174, 184

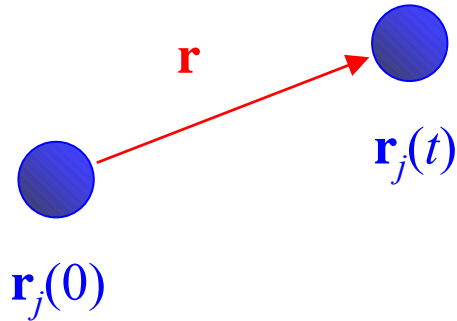
$$D_t = D_o \left. \frac{\partial \ln f}{\partial \ln c} \right|_T$$

sorbate fugacity

$$\lim_{c \rightarrow 0} D_s(c) = \lim_{c \rightarrow 0} D_t(c) = \lim_{c \rightarrow 0} D_o(c)$$



# INCOHERENT QENS SPECTRA FROM MD



Hydrodynamic limit  
(isotropic self-diffusion over large  $r, t$ ):

Self-part of the van Hove correlation function:

$$G_s(\mathbf{r}, t) = \frac{1}{N} \left\langle \sum_{j=1}^N \delta[\mathbf{r} + \mathbf{r}_j(0) - \mathbf{r}_j(t)] \right\rangle$$

$$G_s(r, t) = \frac{1}{(4\pi D_s t)^{3/2}} \exp\left[-\frac{r^2}{4D_s t}\right]$$

Gaussian in  $r$

Self-part of the intermediate scattering function:

$$I_s(\mathbf{Q}, t) = \int G_s(\mathbf{r}, t) \exp(i\mathbf{Q} \cdot \mathbf{r}) d^3 r$$

$$I_s(Q, t) = \exp(-D_s Q^2 t)$$

Gaussian in  $Q$ , exponential in  $t$

Incoherent scattering function:

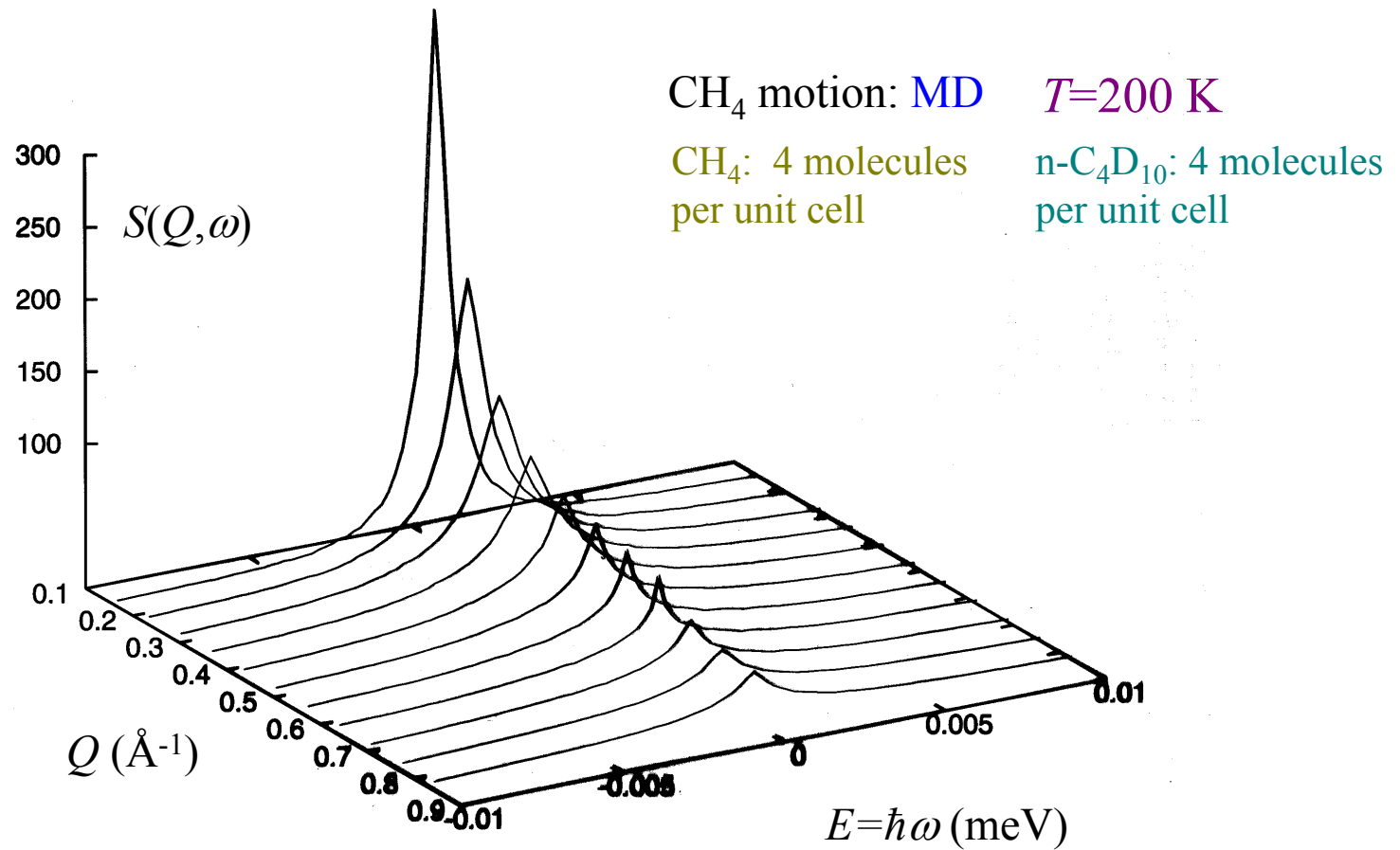
$$S_s(\mathbf{Q}, \omega) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} I_s(\mathbf{Q}, t) \exp(-i\omega t) dt$$

$$S_s(\mathbf{Q}, \omega) = \frac{1}{\pi} \frac{D_s Q^2}{\omega^2 + (D_s Q^2)^2}$$

Lorentzian in  $\omega$

HWHM =  $D_s Q^2$

# DIFFUSION OF METHANE – n-BUTANE MIXTURES IN SILICALITE



L.N. Gergidis and DNT *J. Phys. Chem.*, **103**, 3380-3390 (1999)

L.N. Gergidis, DNT, and H. Jobic *J. Phys. Chem.*, **104**, 5541-5552 (2000)

# SCATTERING FUNCTION LINE SHAPE ANALYSIS: METHANE – n- BUTANE MIXTURES IN SILICALITE

$T=200$  K

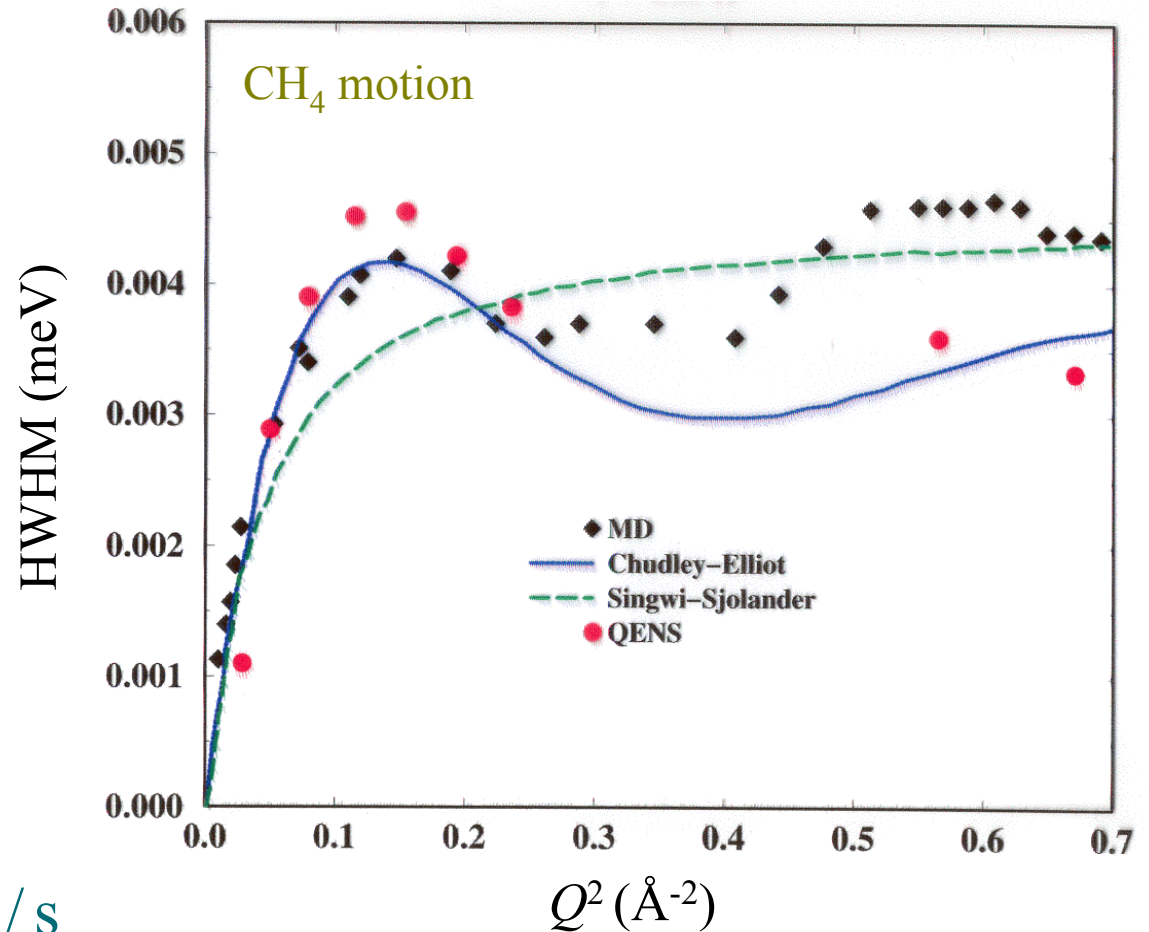
$\text{CH}_4$ : 2 molecules  
per unit cell

$n\text{-C}_4\text{D}_{10}$ : 4 molecules  
per unit cell

MD:

$$D_{s,\text{CH}_4} = 1.8 \times 10^{-5} \text{ cm}^2 / \text{s}$$

$$D_{s,n\text{-C}_4\text{D}_{10}} = 4.1 \times 10^{-6} \text{ cm}^2 / \text{s}$$



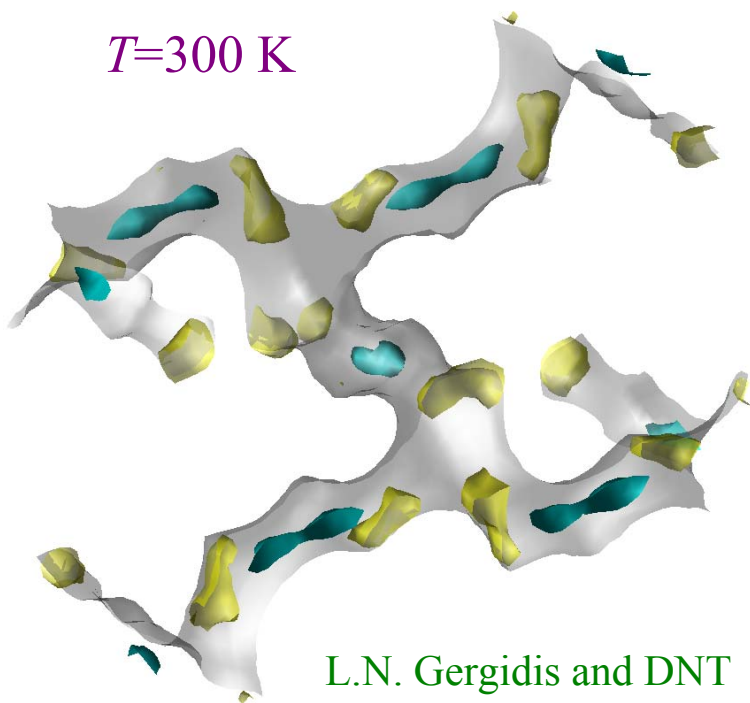
# DIFFUSION OF METHANE – n-BUTANE MIXTURES IN SILICALITE

## Distribution in the pores

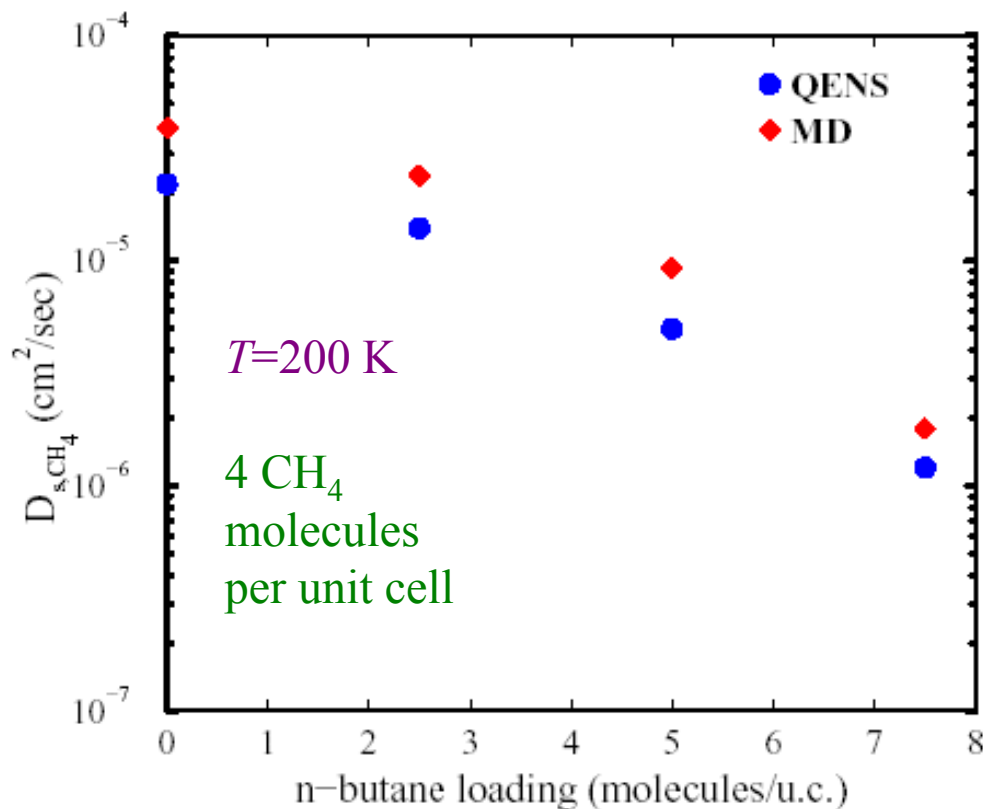
$\text{CH}_4$ : 4  
molecules per  
unit cell

$\text{n-C}_4\text{H}_{10}$ : 9  
molecules per  
unit cell

$T=300$  K



## Methane self-diffusivity



L.N. Gergidis and DNT *J. Phys. Chem.*, **103**, 3380-3390 (1999)

L.N. Gergidis, DNT, and H. Jobic *J. Phys. Chem.*, **104**, 5541-5552 (2000)

# $D_t$ , $D_0$ MEASUREMENT AND SIMULATION

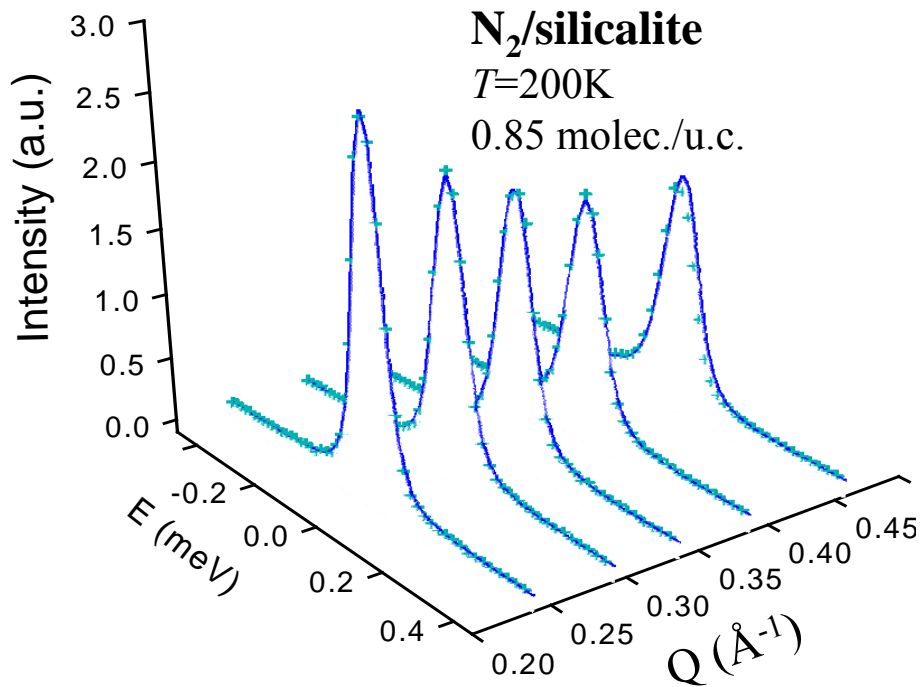
## Coherent Quasielastic Neutron Scattering (QENS)

Dr. Hervé Jobic, Institut de Recherches sur la Catalyse, CNRS, Villeurbanne, France:  
IN6 spectrometer, ILL, Grenoble

Coherent Scattering Function  
(isotropic motion) at low  $Q$ :

$$S_{\text{coh}}(\mathbf{Q}, \omega) = \frac{S(Q)}{\pi} \frac{D_t Q^2}{\omega^2 + (D_t Q^2)^2}$$

$D_t$  extracted from slope of  
HWHM wrt  $Q^2$  at low  $Q$ .



Equilibrium MD Simulation, duration  $\leq 10\text{ns}$

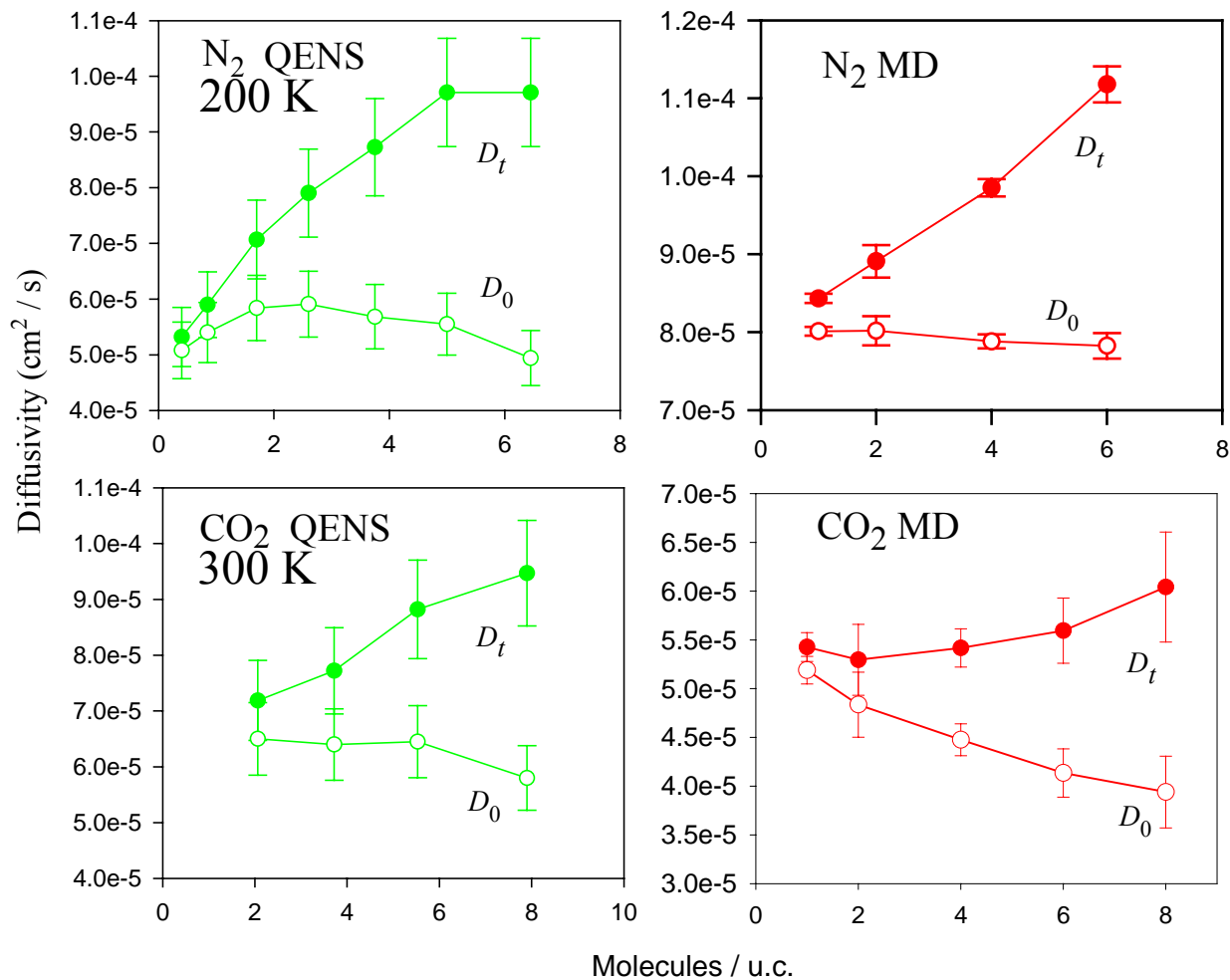
$$D_0 = \frac{1}{3N} \int_0^\infty dt \sum_{i=1}^N \sum_{j=1}^N \langle \mathbf{v}_i(t) \cdot \mathbf{v}_j(0) \rangle = \frac{d}{dt} \lim_{t \rightarrow \infty} \frac{N}{6} \left\langle \left[ \mathbf{R}_{\text{cm}}(t) - \mathbf{R}_{\text{cm}}(0) \right]^2 \right\rangle$$

$\uparrow$  Molecular velocities  
Green-Kubo

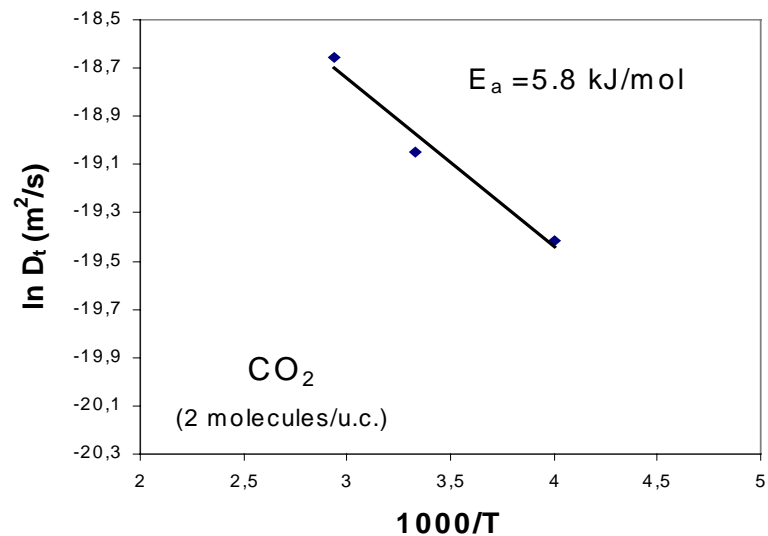
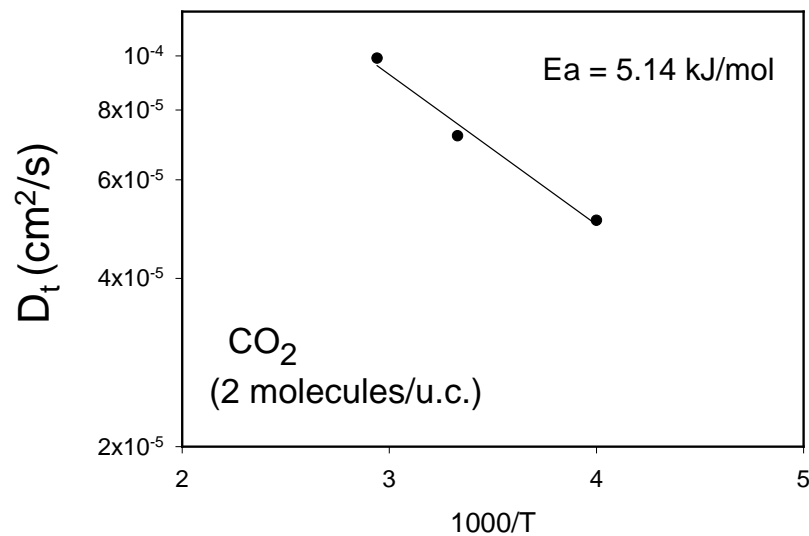
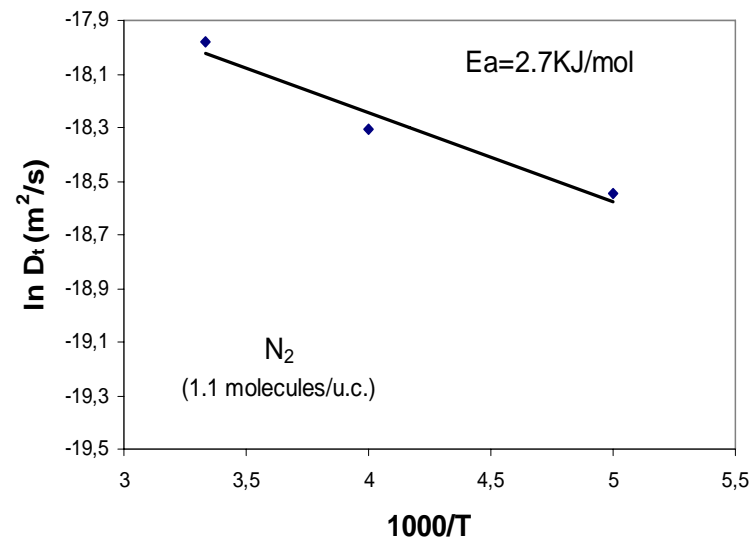
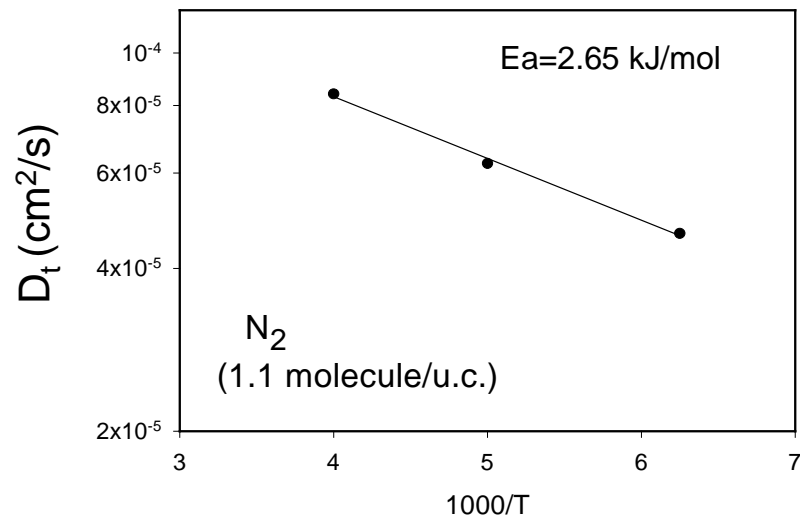
$\uparrow$  Center of mass of swarm of  $N$  molecules: Einstein

# TRANSPORT DIFFUSIVITY OF N<sub>2</sub> AND CO<sub>2</sub> IN SILICALITE-1: COHERENT QENS MEASUREMENTS AND MD SIMULATIONS

G. K. Papadopoulos, H. Jovic, DNT, *J. Phys. Chem. B*, **108**, 12748 (2004)



# TRANSPORT OF N<sub>2</sub> AND CO<sub>2</sub> IN SILICALITE

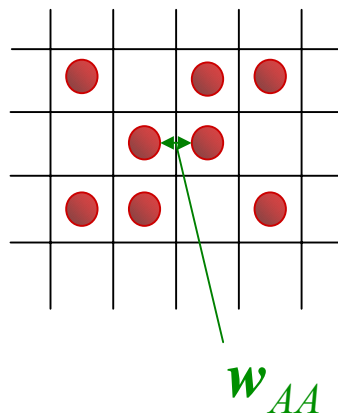


QENS experiment (H. Jobic)

MD (Athens)

# ORIGIN OF DIFFERENT $\theta$ -DEPENDENCE OF $D_0$ OF $N_2$ AND $CO_2$ IN SILICALITE-1

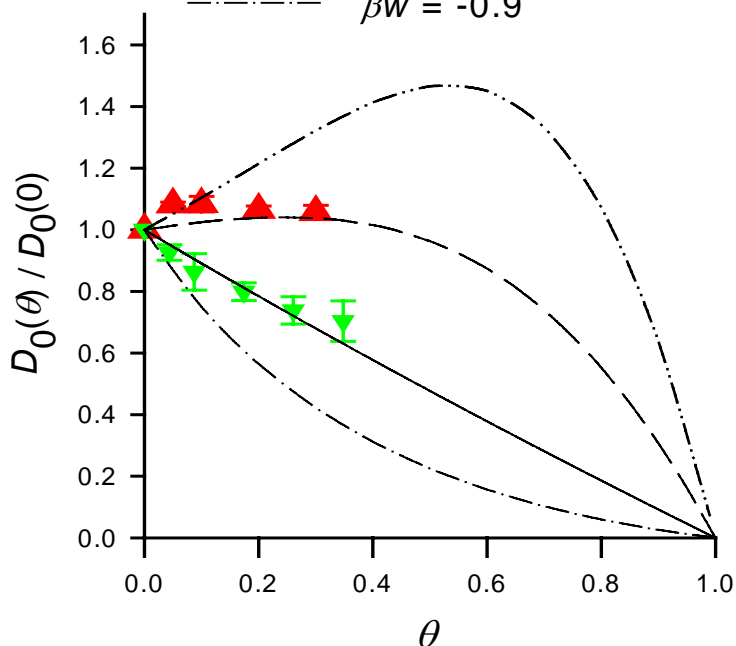
D.A. Reed and G. Ehrlich,  
*Surface Sci.*, **102**, (1981)  
588.



- Regular lattice of sites
- $Z$  coordination number
- $w_{AA}$  nearest-neighbor interactions
- Quasichemical approximation

$$w = \frac{Zw_{AA}}{2}$$

- ▲ MD  $N_2$
- ▼ MD  $CO_2$
- $\beta w = -0.1$
- - -  $\beta w = 0.7$
- · - ·  $\beta w = 1.15$
- - -  $\beta w = -0.9$



**Molecular jumps with fixed attempt rate. Attempt to jump into an occupied site unsuccessful.**

Decreasing trend of  $D_0(\theta)$  for  $CO_2$  explained in terms of stronger, attractive sorbate-sorbate interactions in that system.

G.K. Papadopoulos, H. Jobic, DNT, *J.Phys.Chem. B*, **108**, 12748 (2004).



# RATIOS OF PURE GAS PERMEABILITIES THROUGH MFI MEMBRANES AT 300 K

**Experiments:** Supported Membrane Permeation

Y. Yan, M.E. Davis, G.R. Gavalas, *Ind. Eng. Chem. Res.* **34**, 1652 (1995).

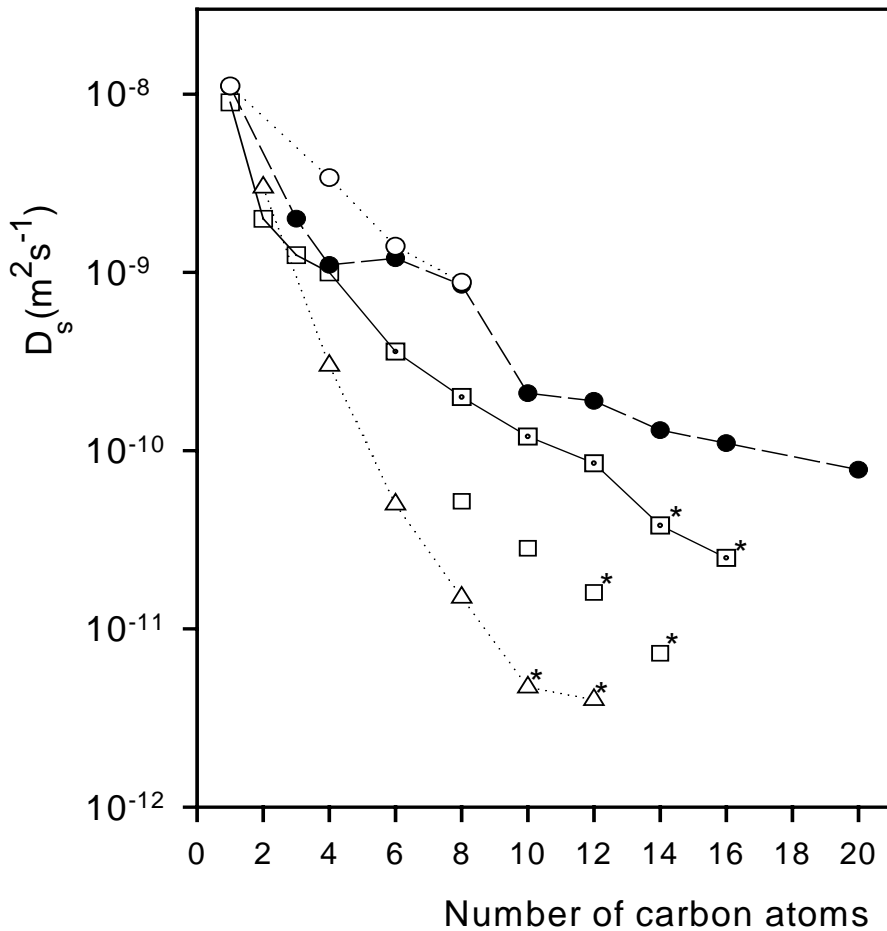
**Simulations:** Equilibrium NVE MD and GCMC

K. Makrodimitris, G.K. Papadopoulos, DNT, *J. Phys. Chem. B* **777**, 105 (2001).

Sorbates	Experiments	Simulations
CH <sub>4</sub> /N <sub>2</sub>	1.2	1.6
CO <sub>2</sub> /CH <sub>4</sub>	2.3	2.2
CO <sub>2</sub> /N <sub>2</sub>	2.8	3.4

# SELF-DIFFUSIVITY OF *n*-ALKANES IN SILICALITE-1

$T=300$  K, low occupancy



(○) MD simulations [F. Leroy, B. Rousseau, and A.H. Fuchs, *Phys. Chem. Chem. Phys.* 6 (2004) 775 ]

(●) hierarchical simulations [E.J. Maginn, A.T. Bell, and DNT, *J.Phys.Chem.* 100 (1996) 7155 ]

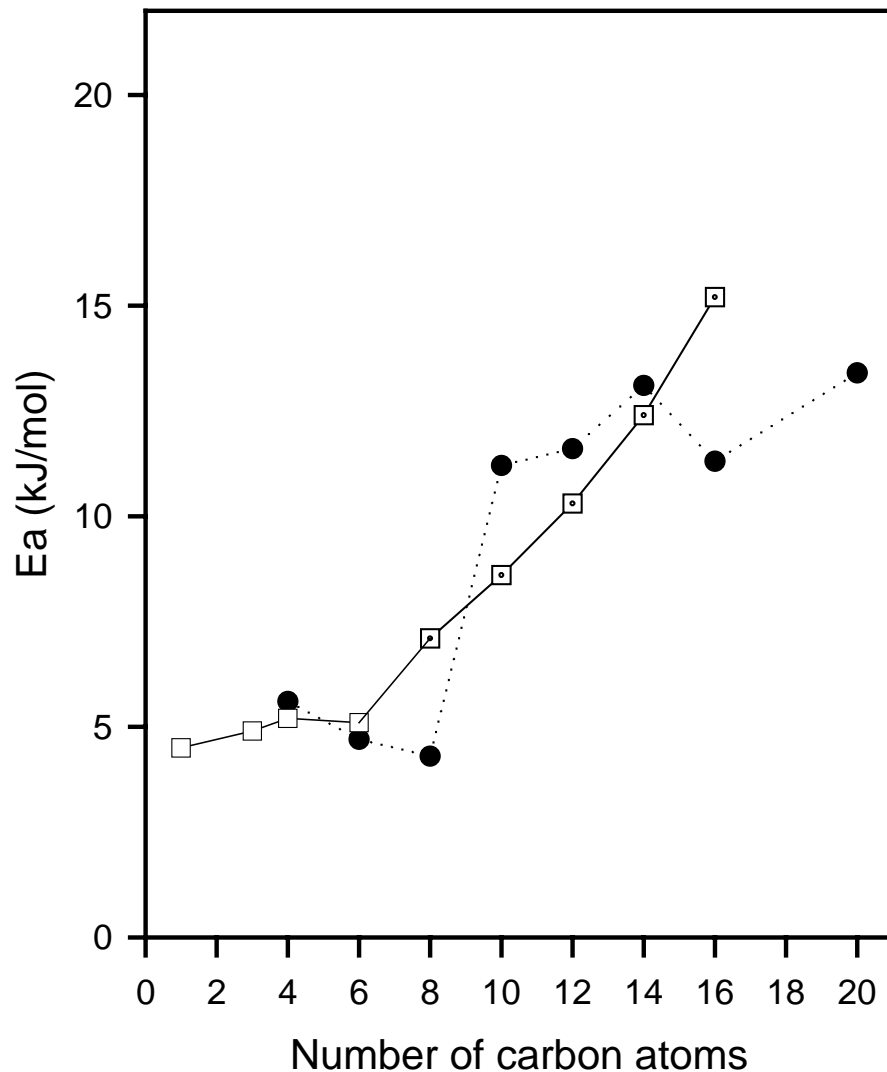
(□) QENS in Na-ZSM-5 [H. Jobic, *J. Mol. Catal. A-Chem.* 158 (2000) 135-142]

(□) QENS in silicalite-1 [H. Jobic and DNT, *J. Phys. Chem.* 110 (2006) 1964].

(△) PFG NMR [H. Jobic, W. Schmidt, C.B. Krause, and J. Kärger, *J. Micropor. Mesopor. Mater.* 90, 299(2006)].

Asterisked symbols: extrapolation to 300 K.

# ACTIVATION ENERGY FOR SELF-DIFFUSION: ALKANES IN SILICALITE-1



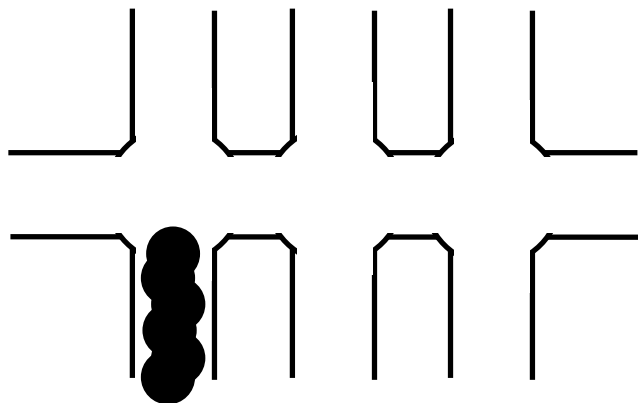
low occupancy

(●) hierarchical simulations [E.J. Maginn, A.T. Bell, and DNT, *J.Phys.Chem.* 100 (1996) 7155 ]

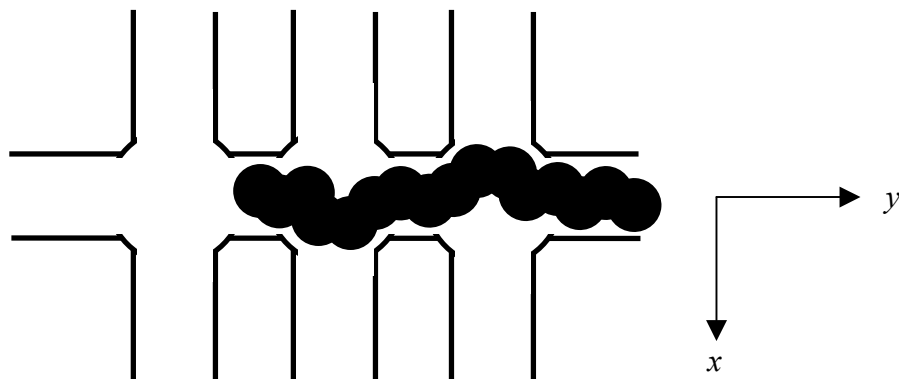
(□) QENS in ZSM-5 [H. Jobic, *J. Mol. Catal. A-Chem.* 158 (2000) 135-142 ]

(□) QENS in silicalite-1 [ H. Jobic and DNT *J. Phys. Chem.* 110 (2006) 1964 ].

## POSSIBLE EXPLANATION OF CHAIN LENGTH-DEPENDENCE OF $E_a$



*n*-hexane sorbed in a Z  
coarse-grained conformation  
in silicalite-1



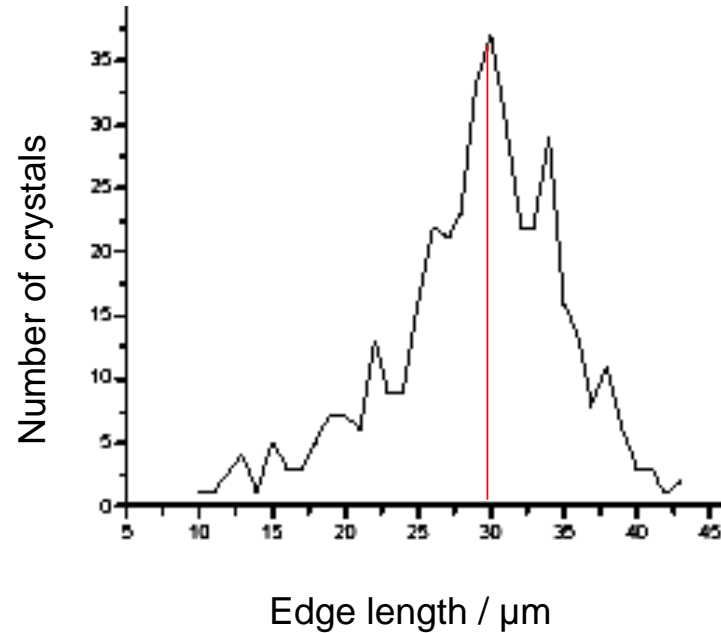
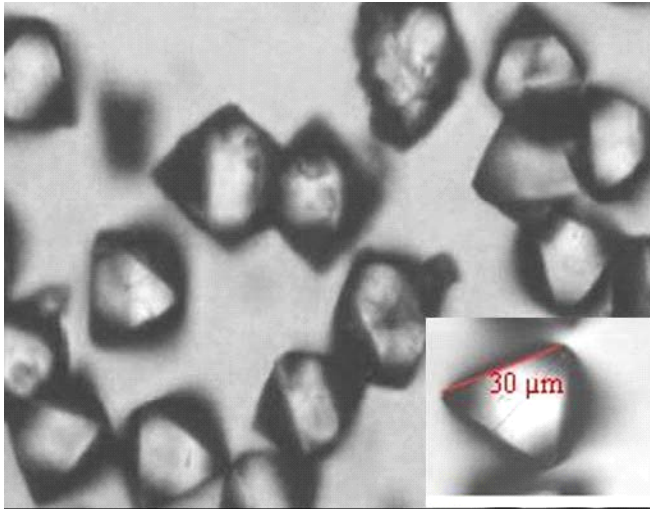
*n*-hexadecane sorbed in a SISIS  
coarse-grained conformation  
within silicalite-1

$n < 8$ : Jumps across more spacious, energetically less favorable intersection regions

$n > 8$ : Conformational rearrangement accompanies translation along S channels

# PARTICLE-BASED COMPUTER RECONSTRUCTION

## Experimental input



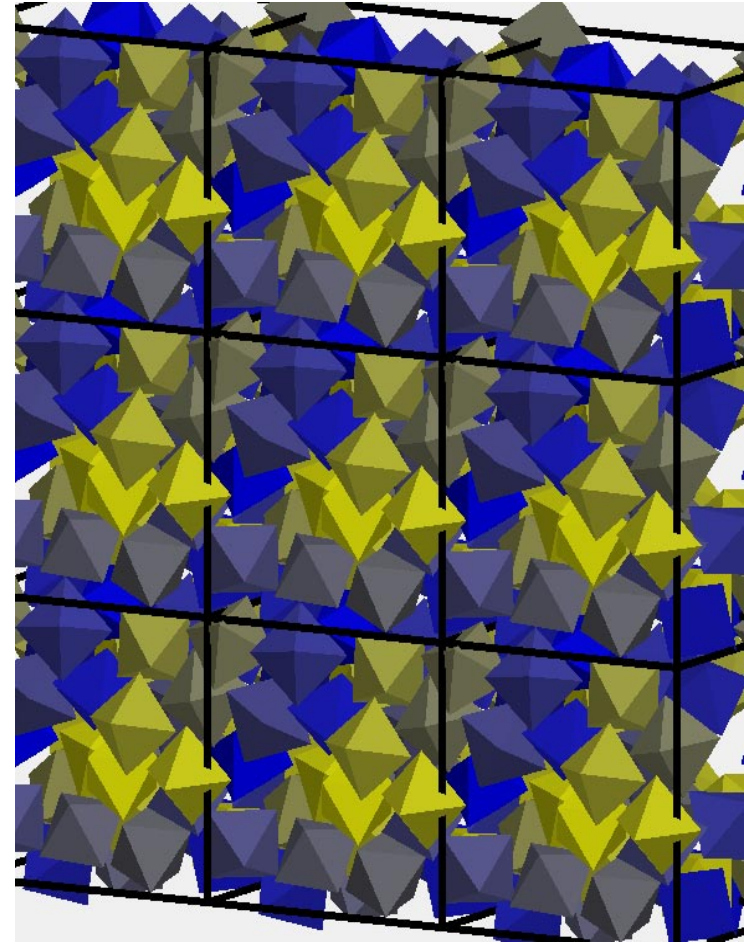
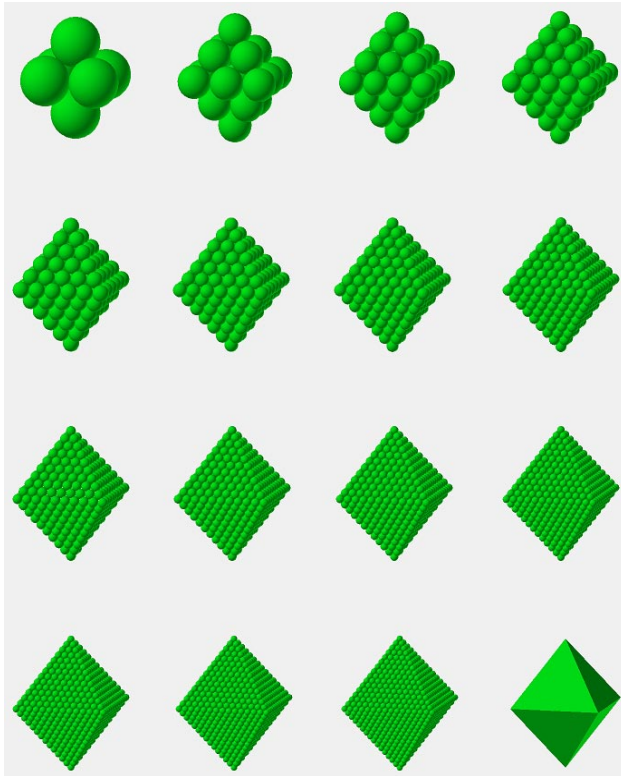
$$\rho_{\text{crystal}} = 1530 \text{ kg / m}^3$$

$$\rho_{\text{bed}} = 590 \text{ kg / m}^3$$

**Bed of NaX zeolite crystals:** Dr. S.Vasenkov , Prof. Dr. J. Kärger, U. Leipzig

# PARTICLE-BASED COMPUTER RECONSTRUCTION

Application: bed of NaX crystals



Sequence of energy minimizations at fixed porosity, using progressively higher generations for the representation of particles.



# SIMULATION OF “LONG-RANGE” DIFFUSIONAL PROCESSES IN INTERCRYSTALLINE SPACE

Track long trajectories of individual sorbate molecules by kinetic MC simulation:

- Translational motion with constant speed  $\bar{u}$
- Intermolecular collisions with mean free path  $\lambda$
- Diffuse reflections upon collision with solid surfaces

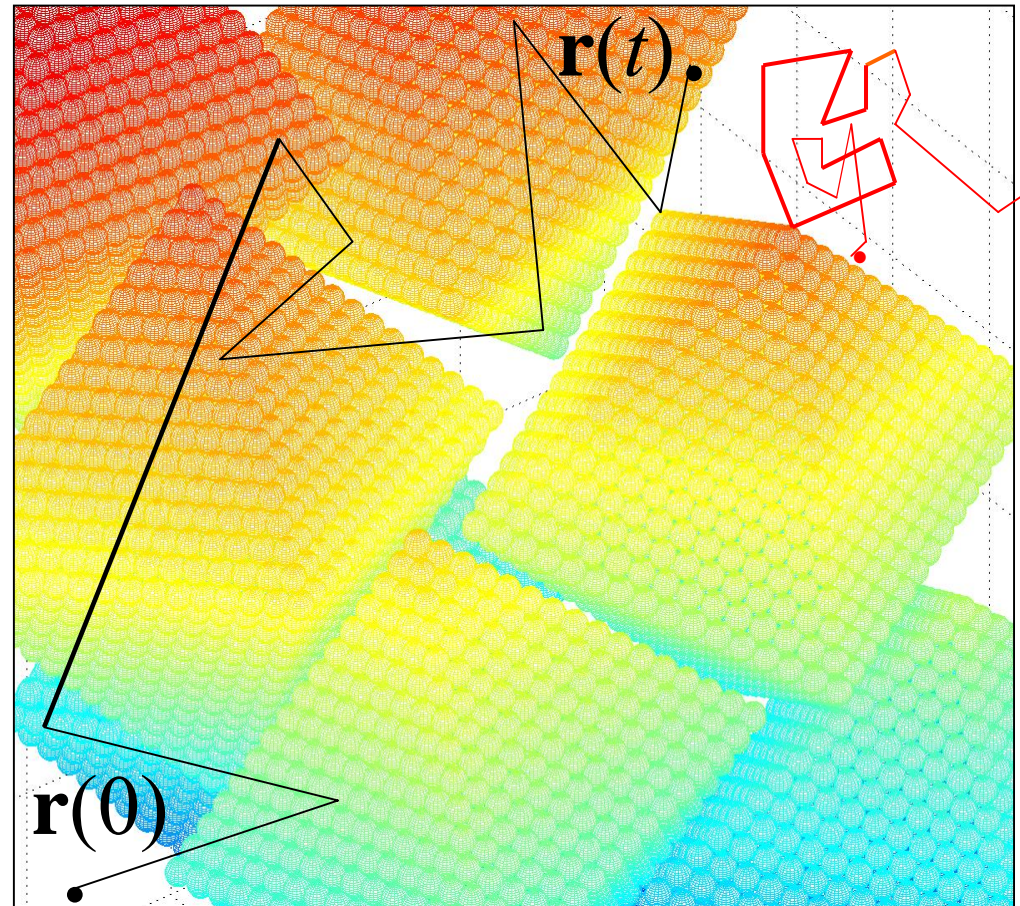
$$\lambda = \frac{k_B T}{\sqrt{2\pi}\sigma^2 P}$$

$$\bar{u} = \sqrt{\frac{8k_B T}{\pi m}}$$

Effective Diffusivity:

$$D = \lim_{t \rightarrow \infty} \frac{\langle [\mathbf{r}(t) - \mathbf{r}(0)]^2 \rangle}{6t}$$

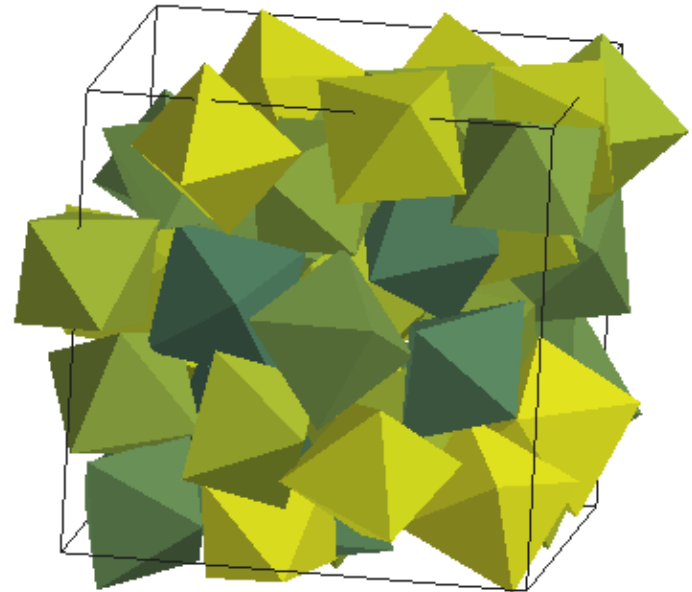
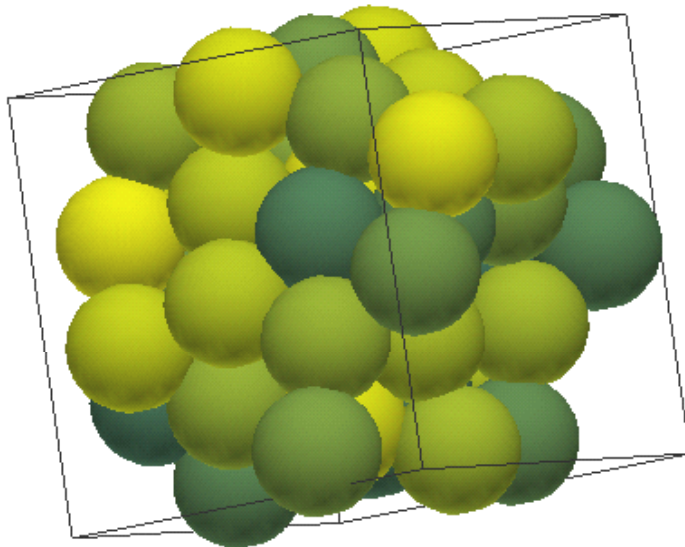
Measurements: O. Geier, S. Vasenkov, J. Kärger, *J. Chem. Phys.* **2002**, *117*, 1935-1938



# LONG-RANGE DIFFUSION IN CRYSTAL BEDS

Ethane in two model beds:

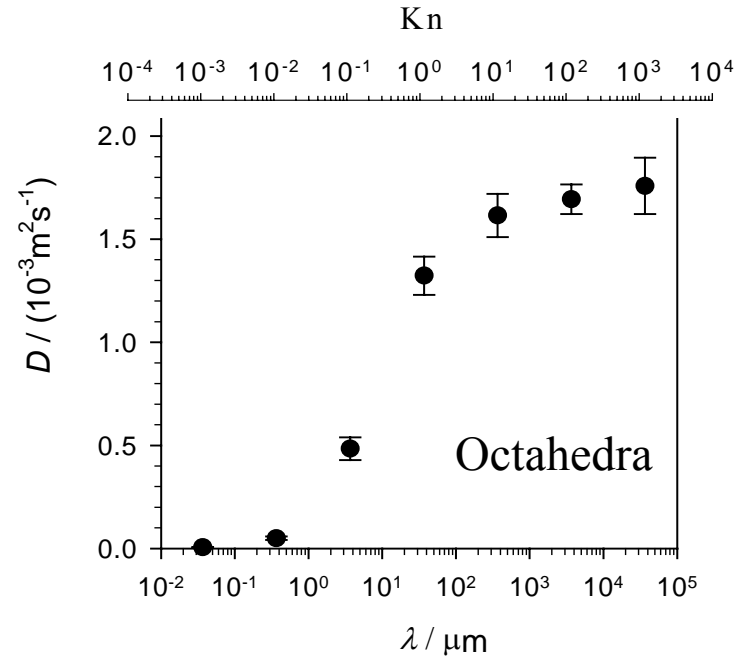
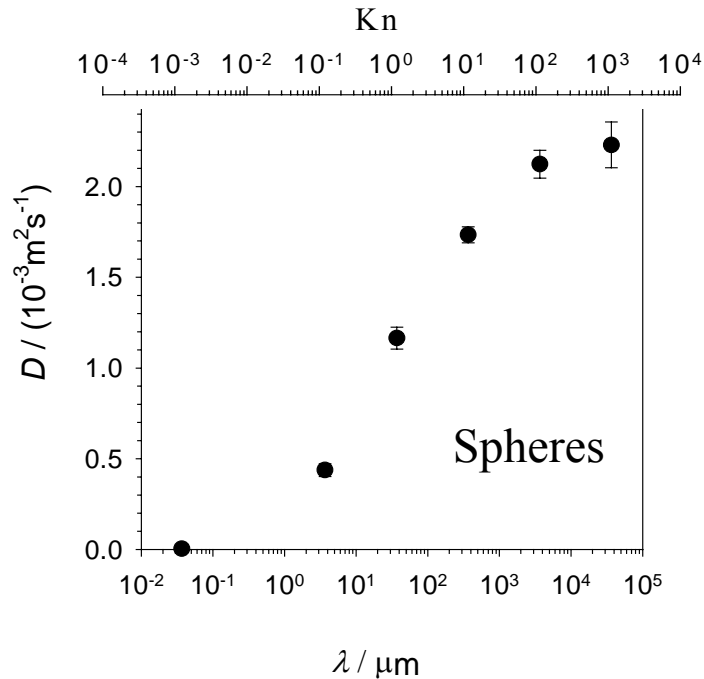
Crystal geometry	Simulation box length ( $\mu\text{m}$ )	Particle size ( $\mu\text{m}$ )	$\phi$	Generation
Spherical	116.740	28.969 (diameter)	0.60	0
Octahedral	116.740	29.996 (edge)	0.60	15





# TRANSITION FROM MOLECULAR TO KNUDSEN DIFFUSION WITH INCREASING $\lambda$

Ethane in NaX bed, 295 K,  $\phi = 0.60$



Molecular diffusion regime

$$D = D_b = \frac{D_{0b}\phi}{\eta_b} \leftarrow \text{tortuosity factor}$$

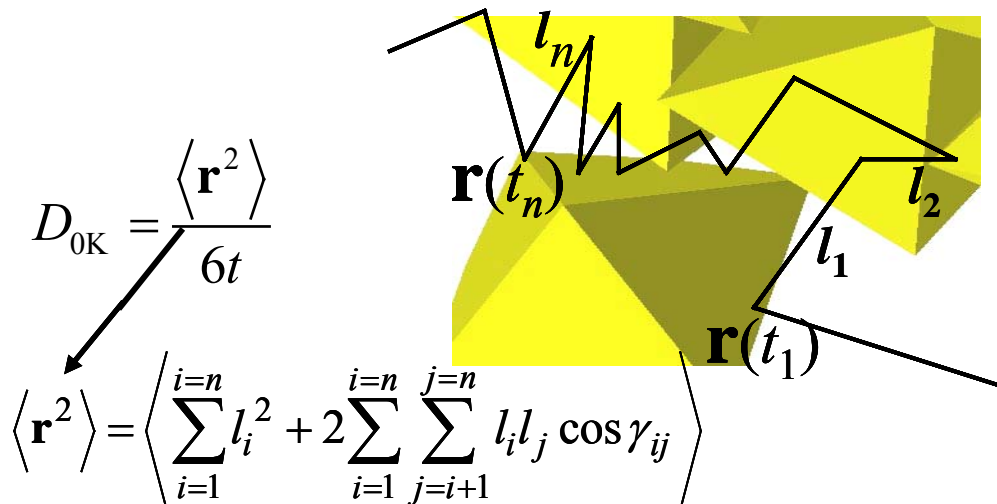
$$D_{0b} = \frac{1}{3} \bar{u} \lambda$$

Knudsen regime

$$D = D_K = \frac{D_{0K}\phi}{\eta_K}$$

$$D_{0K} = \frac{1}{3} \bar{u} l_{ss} \leftarrow \text{mean solid-to-solid path length}$$

# DERJAGUIN CORRECTION TO THE KNUDSEN DIFFUSIVITY



$$D_{0K} = \frac{\langle \mathbf{r}^2 \rangle}{6t}$$

$$\langle \mathbf{r}^2 \rangle = \left\langle \sum_{i=1}^{i=n} l_i^2 + 2 \sum_{i=1}^{i=n} \sum_{j=i+1}^{j=n} l_i l_j \cos \gamma_{ij} \right\rangle$$

Consecutive lengths not correlated  
Lengths and orientations mutually independent

$$\langle \mathbf{r}^2 \rangle = n \langle l^2 \rangle + 2n \langle l \rangle^2 \sum_{k=1}^{\infty} \langle \cos \gamma_k \rangle$$

$$D_{0K} = \frac{1}{3} \bar{u} \langle l \rangle \left[ \frac{\langle l^2 \rangle}{2 \langle l \rangle^2} - \beta \right]$$

$l_{SS}$   
↓

$$\beta = - \sum_{k=1}^{\infty} \langle \cos \gamma_k \rangle$$

# TORTUOSITY FACTORS FROM EXPERIMENT AND SIMULATION

$\phi$	$d_s$ ( $\mu\text{m}$ )	$l_{ss}$ ( $\mu\text{m}$ )	$\eta_K/\eta_b$ PFG NMR*	$\eta_K/\eta_b$ simulation	$\beta$	$\frac{\langle l^2 \rangle}{2\langle l \rangle^2}$	$\eta_K/\eta_b$ corrected
0.60 (sph)	26.51	26.35	-	1.55	0.309	0.935	0.97
0.60 (oct)	24.49	24.32	4.8-10	2.04	0.298	0.849	1.12

$$d_s = \frac{4V_{\text{void}}}{S}$$

$l_{ss}$  = mean solid-to-solid path length

\*O. Geier, S. Vasenkov, J. Kärger, *J. Chem. Phys.* **2002**, *117*, 1935-1938

# CONCLUSIONS

- Molecular simulations are useful in elucidating mechanisms and predicting rates of diffusion in zeolitic materials.
- Molecular simulations aid in the interpretation of diffusion measurements.
- Broad spectra of length and time scales present in technologically important zeolitic materials necessitate the development of hierarchical modelling and simulation approaches.

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